**Constraining Electron Parallel Energy in Electrostatic Fields through the Anomalous Doppler Effect Induced by External Electromagnetic Waves**

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The interaction between free electrons and electromagnetic waves under the influence of magnetic and electrostatic fields is investigated using a Volume-Preserving algorithm. When the electric field of the EMW, containing a left-hand polarization component, exceeds a critical threshold, it facilitates continuous transfer of parallel electron energy into rotational energy through the Anomalous Doppler Effect (ADE). This process transforms the electric field's work along the magnetic field into perpendicular kinetic energy, leading to saturation of the electron’s parallel kinetic energy and continuous growth of its perpendicular kinetic energy. A theoretical model based on energy, momentum, and angular momentum conservation elucidates the role of left-hand polarization in the Anomalous Doppler Effect and provides a generalized framework for interpreting electron-wave interactions. This study proposes a novel approach for mitigating runaway electrons in magnetically confined plasmas, suggesting the use of extraordinary waves launched from the high-field side with an energy flux of several watts per square meter to saturate parallel energy in Tokamaks.

**I. Introduction**

In the beginning of burning plasma device discharge (current ramp up phase), the magnetohydrodynamic (MHD) instabilities and disruption can generate quasi-static toroidal electric fields that accelerate electrons to energies reaching several tens of MeV. This acceleration occurs when the force exerted by the quasi-static electric field surpasses the opposing forces from radiation and collisional drag. These high-energy electrons, known as runaway electrons, can inflict severe damage on the tokamak’s interior walls, thereby shortening the device’s operational lifespan. An intriguing possibility is to convert the energy gained by electrons from quasi-static electric fields into rotational energy within the magnetic field. This approach not only suppresses the energy of runaway electrons, reducing their harmful impact on the device, but also improves discharge performance by minimizing the consumption of ohmic field energy.

The transport of parallel energy from electrons into rotational energy primarily occurs through three different mechanisms, including (1) the electron avalanche process [25], (2) collision less pitch-angle scattering [22], and (3) the Anomalous Doppler Effect [15]. Current strategies to suppress runaway electrons, such as gas injection [34] and the enhancement of magnetic perturbations [36], often have unintended side effects and disrupt the discharge environment. In contrast, the Anomalous Doppler Effect provides a cleaner mechanism, making it a particularly attractive avenue for further investigation.

When electrons move in static magnetic fields and interact with external electromagnetic waves (EMW) of frequency ω and wave vector k⃗, they undergo a scattering phenomenon under the resonant condition, where and .This scattering results in the transfer of momentum from parallel motion to rotational motion, a phenomenon known as the Anomalous Doppler Effect. The Anomalous Doppler Effect was first thoroughly described in the seminal works of Ginzburg and Frank [15, 10, 24].

Recently, the Anomalous Doppler Effect has garnered increasing attention in fields such as space radiation [8], runaway electron instabilities [21], and materials science [27]. It is believed that Anomalous Doppler Effect can explain phenomena like whistler turbulence in solar flare loops [8], the step-like structure in Electron Cyclotron Emission (ECE) observed in tokamaks [23, 6, 5], and the microwave bursts during Edge Localized Modes (ELMs) [12]. Furthermore, Anomalous Doppler Effect has shown potential for suppressing runaway electron energy in tokamak discharges.

This potential was demonstrated by F. Santini [26], who found that high-energy runaway electrons could be significantly reduced through Anomalous Doppler Effect during lower hybrid wave heating in the Frascati Tokamak . However, it is important to note that the high power of lower hybrid waves also increases the population of nonthermal electrons through Landau resonance, leading to a subsequent rise in runaway electrons after the lower hybrid waves are turned off. This side effect poses a challenge to the use of lower hybrid waves for suppressing runaway electrons.

The Anomalous Doppler Effect has been explored in various experiments and theories [7, 4, 17, 28, 30, 13, 19, 14, 18], but previous discussions of the single-particle Anomalous Doppler Effect have mainly focused on either quantum or classical theories, without considering the role of the accelerating electrostatic field [16, 26, 3]. Understanding the Anomalous Doppler Effect in the presence of electrostatic fields is essential for comprehending the physics of pitch-angle scattering of runaway electrons by electromagnetic waves in Tokamak discharges. Additionally, the behavior of Anomalous Doppler Effect in classical electrodynamics remains poorly understood, primarily due to the complexity of the formulas in earlier classical theories. It is still unclear why parallel kinetic energy can convert into transverse internal energy during Anomalous Doppler Effect resonance, and which types of electromagnetic waves can trigger it.

This paper presents a direct simulation of full orbit electron motion in uniform magnetic fields, along with accelerating electrostatic and electromagnetic fields, using the Volume-Preserving Algorithm [37]. Compared to conventional algorithms like Boris [38], the Volume-Preserving Algorithm ensures long-term accuracy and conservativeness through a systematic splitting method, making it an ideal approach for nonlinear electron dynamic simulations. To directly observe the Anomalous Doppler Effect, an electron is placed in a uniform magnetic field and an electrostatic field, which is oriented opposite to background magnetic field. This setup allows the electron to be accelerated parallel to background magnetic field. During the simulation, a slow electromagnetic wave with a phase velocity smaller than that of light in vacuum is introduced as an induced wave. This wave enables us to observe the effects when the electron’s velocity reaches the resonant condition for the Anomalous Doppler Effect. We explore resonance with three types of polarization waves: linear polarization, left-hand circular polarization, and right-hand circular polarization. The results show that only the wave with left-hand circular polarization induces the Anomalous Doppler Effect for runaway electrons. The simulation also reveals the critical energy of waves at which the electron's parallel velocity is constrained and consistently transfers parallel energy from the electrostatic field to transverse rotational energy. Furthermore, the self-consistency between quantum theory and direct simulation of the Anomalous Doppler Effect is examined. The analysis of dispersion, polarization, and resonant moments, we determine that the extraordinary wave is most suitable for triggering the Anomalous Doppler Effect in plasma. Based on these findings, we propose an effective method for controlling runaway electrons.

Section II discusses the quantum theory. The numerical simulation framework and results are presented in Section III. The trapping threshold is examined in Section IV. Section V explores the dynamics of electromagnetic waves driving the Anomalous Doppler Effect in magnetized plasma. The runaway electron suppression method using extraordinary wave injection is introduced in Section VI. Finally, the summary is provided in Section VII.

**II. Quantum Theory of the Anomalous Doppler Effect**

This extraordinary phenomenon has been previously discussed in terms of energy conservation by V.L. Ginzburg [14], I. Tamm [31], Nezlin [24], and I.M. Frank [11]. In this work, we provide an analysis based on the conservation of angular momentum. As illustrated in Fig. 1, when charged particles move through a medium at speeds greater than the speed of light in that medium, induced currents are generated. These currents, in turn, stimulate secondary waves that interfere with the electromagnetic field of the moving particles, resulting in Cherenkov radiation. The direction of Cherenkov radiation is constrained to the Cherenkov radiation angle , where c′ is the speed of light in the medium and v is the velocity of the charged particles.

In this case, the charged particle is replaced with a system that has internal energy, such as an oscillator or a cyclotron electron in a magnetic field. When the system moves faster than the speed of light , it emits photons with angular frequency and wavevector in the direction . The direction of the emitted photon is not influenced by the interference of secondary waves and can occur in any direction, as shown in Fig. 2.

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Figure 1. Schematic diagram of Cherenkov Radiation. The red points stand for the snapshot of the electron at different time.

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Figure 2. The blue region shows the Anomalous Doppler Effect.

According to energy conservation and momentum conservation:

(1)

(2)

In the above, and represent the kinetic energy and internal energy of the system before emitting a photon, and and represent the energy of the system after emitting a photon. With the assumption that photons energy is far less than the initial kinetic energy , the losses of kinetic energy after emitting a photon can be expressed as , where is the velocity of the system before emitting a photon, and also .

(3)

Here, . While the system velocity is greater than the speed of light in the medium . According to the sign of , we can divide radiation into three regions, as shown in Fig. 2.

1. For , . The system produces photons by consuming its own internal and kinetic energy, this region refers to the Normal Doppler Effect (NDE).
2. For , , the loss of kinetic energy by the system is completely converted into photon energy; this line refers to the Cerenkov Effect.
3. For , , this region is referred to the Anomalous Doppler Effect (ADE), where the system gains internal energy after emitting photons. It means the loss of kinetic energy is converted to photons and the system’s internal energy.

All three effects are possible when the system velocity exceeds the speed of light . While the system velocity is less than the speed of light , only Normal Doppler Effect exists. As observed, the type of phenomenon can be determined by examining the change in internal energy after the emission of photons.

The freely moving electron has a velocity along the background magnetic field and a velocity perpendicular to the magnetic field, as shown in Fig. 3. The kinetic energy is . The referred to is the Lorentz factor The internal energy represented as . According to the angular momentum conservation, we have the Eq.4.

(4)

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Figure 3. Schematic diagram of electron cyclotron system. , .

The emitted photon is considered to contain angular momentum of , and the angular momentum of the cyclotron electron before and after emitting the photon is and , respectively. Since the magnetic field is aligned along the z-direction, the angular momentum should be equal to .

According to the quantum theory, the electron wave in the static magnetic field can be presentated as Eq. 5.

(5)

With the term is a normalized coefficient, A is the vector potential and s is the position.

The intrinsic equation of angular momentum in the z-direction has been transformed as Eq. 6.

(6)

(7)

With ,and , the Eq.7 is presentated as Eq. 8.

(8)

With is the electron rest mass, is the Lorentz factor and is the electron cyclotron frequency in rest frame.

The variation in the angular momentum of the electron along z is presented as Eq. 9.

(9)

With m is the number of photon’s angular momentum in z direction.

The internal energy changes , with the Eq. 9, will be transformed as Eq. 10.

(10)

According to the Eq. 3 and Eq. 10, the photon could be presented as Eq. 11 and Eq. 12.

(11)

(12)

Here, represents the loss of kinetic energy , represents the energy of the photon, and represents the change in the electron cyclotron energy (internal energy change). There are three scenarios about the the internal energy changes.

1. With , , the cyclotron electron internal energy decreases after emitting a photon, and the emitted photon will have right-hand circular polarization with angular momentum to maintain angular momentum conservation. This process is called the Normal Doppler Effect .
2. For , , the Cherenkov Effect occurs, where the emitted photon does not cause any change in the internal energy of the cyclotron electron.
3. With , , the Anomalous Doppler Effect (ADE) occurs,resulting in an increase in the internal energy of the cyclotron electron and the emission of left-hand circular polarization with angular momentum .

As a result, the resonant condition is strongly associated with the wave’s angular momentum. For a plane wave, the wave angular moment number includes only . While for , it indicates that the resonant wave possesses a helicon structure. Based on the above discussion,there are three kinds of resonance for a system when electron moves along the uniform magnetic field with velocity under external EMW ;the resonant frequencies are Normal Doppler frequency, Cerenkov frequency , and Anomalous Doppler frequency. We only include the dominate resonance .

Normal Doppler Effect frequency:

(13)

Cerenkov Effect frequency:

(14)

Anoumalous Doppler Effect frequency:

(15)

With is the angle between background mangetic field and wavevector . These equations are quite common resonant conditions for the kinetic equation of plasma ,what is intriguing is how the motion of electrons differs under various resonant conditions with an electrostatic field.

**III. Numerical Simulation Framework and Result Discussion**

The uniform magnetic field is set on the z-direction. The electron is accelerated by the electrostic field , which on the opposite direction to as shown in Fig. 4. For the dynamics analysis of electrons during interactions with an electromagnetic field, a plane electromagnetic wave is estblished, which characterized by frequency and wavevector .

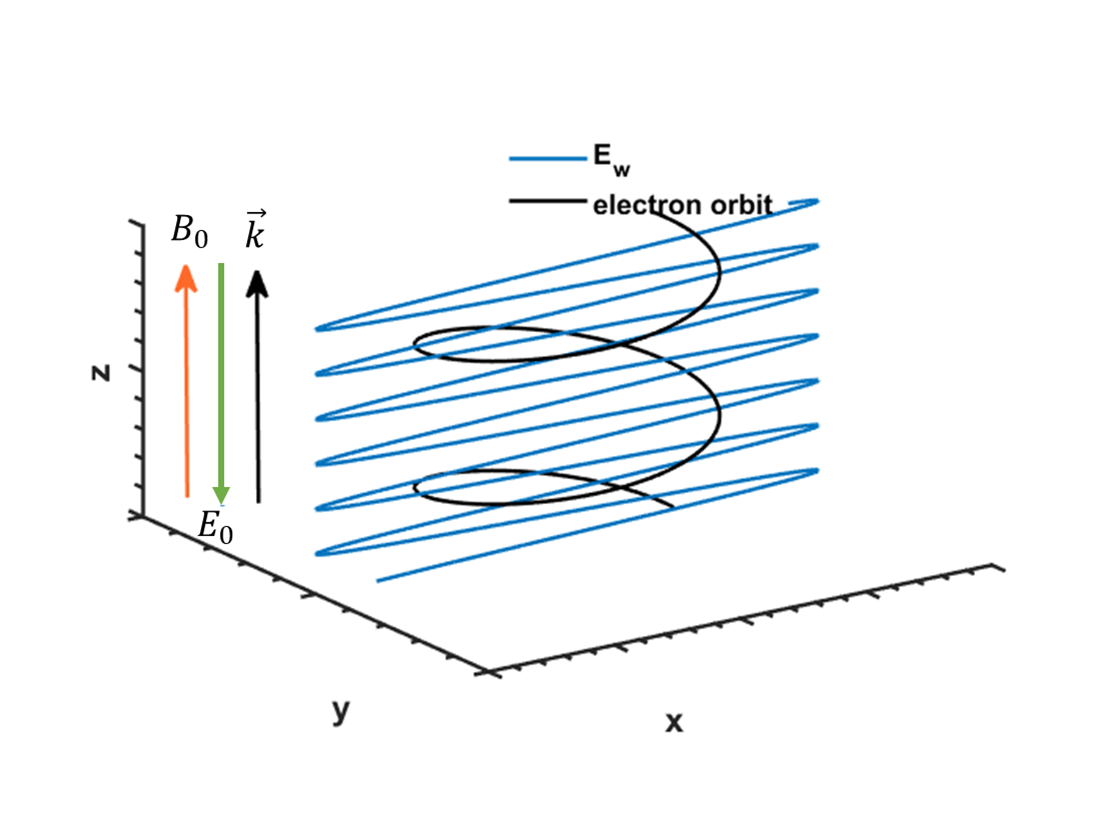


Figure 4. The uniform background magnetic is set on z direction (orange). The electrostatic field is marked with green. The electromagnetic field progates along z direction, with the linear polarization along x direction. The electron orbit has been plot in black.

The electron orbit and motion p in this scenario is presented as Eq. 16. The and are the total field contains static field and electromagnetic field.

(16)

The discrete stucture of Eq.16 is rewrited as Eq. 17 by employing the Volume-Preserving Algorithm[32,33, 37]. The operator Cay(A) denotes the Cayley transform of matrix A[38].

(17)

The dimensionless magnetic matrix is presented as Eq. 18.

(18)

The dimensionless parameters are momentum , magnetic field , total electric field , time step , and position respectively, where the is the electron cyclotron period and .

As a preliminary validation calculation, the parameters are set as follows. The background magnetic field = T . The wave angle frequency is , where . The wavevector . The amplitude of the electric field of the electromagnetic wave is , and the electrostatic field is . All these parameters are only set for the purpose of rapid simulation. The real tokamak scale calculation will be discussed in the following section. The time step is always chosen to satisfy to ensure the accuracy of the simulation. The electron begins at rest and gradually gains speed. The resonant frequency increases according to Eq. 13, 14, and 15.

A diagram of a coil with a wire

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Figure 5. (left) Orbit trajectory of electron motion. (right) Momentum phase space of electron motion.

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Figure 6. Kinetic evolution of electrons in a magnetic field with electromagnetic wave during acceleration. (a) Wave frequencies of anomalous Doppler frequency, normal Doppler frequency, and source wave frequency. (b) The translational velocity differnce in the case with and without the electromagnetic wave. (c) The change of parallel velocity caused by the electromagnetic wave. (d) Cyclotron velocity . (e) Interaction with Linear ,Right-hand circular ,Left-hand circular polarization .

Figure 5 illustrates the evolution of the electron’s orbit and velocity phase during acceleration. The details of the electron’s motion are shown in Figure 6. As the electron accelerates in the electrostatic field (Fig. 6b), the resonant frequencies increase concurrently (Fig. 6a). At around , when the normal Doppler frequency matches that of the induced wave, the perpendicular velocity (or rotational velocity) increases rapidly (Fig. 6d). The parallel velocity induced by the electromagnetic wave also increases, as shown in Fig. 6c. This change can be calculated as , where is the parallel velocity due to both the electromagnetic wave and the electrostatic field, and is the parallel velocity resulting only from the electrostatic field.

This phenomenon corresponds to the Normal Doppler Effect, where the resonant velocity is "subluminal." The absorption of induced waves by the cyclotron electron results in an increase in both parallel and perpendicular velocities, which can be considered a reverse process to the photon emission described earlier. The the Normal Doppler Effect process is widely used for current drive [2] and plasma heating [20] in tokamaks. However, it is generally believed that current drive via electromagnetic waves follows the Fisch mechanism [9], due to the limited toroidal momentum injected by the waves.

The resonant condition is quickly disrupted as the parallel velocity continues to increase until it reaches , at which point the Anomalous Doppler Effect begins to emerge. When the time reaches , the system starts resonating with the induced wave through the Anomalous Doppler Effect, where as shown in Fig. 6a. At this point, the parallel velocity begins to scatter into the perpendicular direction, evident from the decrease in and the increase in as seen in Fig. 5c and Fig. 6d. The resonant condition rapidly disappears as the parallel velocity exceeds the resonant region. To determine which type of electromagnetic wave is responsible for the the Normal Doppler Effect, and the Anomalous Doppler Effect separately, we decompose the linearly polarized wave into left-handed and right-handed circular polarizations. We observe that the right-hand circular polarized wave is responsible for the Normal Doppler Effect, while the left-hand circular polarized wave induces the Anomalous Doppler Effect, as shown in Fig. 6e, which aligns well with our previous analysis.

This phenomenon is understood through the conservation of angular momentum and linear momentum. Electrons exhibit right-handed spin orbital angular momentum in a magnetic field. When an electron absorbs right-handed electromagnetic waves propagating in the parallel direction, the conservation of momentum and angular momentum causes an increase in both the parallel momentum and rotational energy of the electron, corresponding to the Normal Doppler Effect. Conversely, when the electron emits left-handed circular polarization electromagnetic waves propagating in the parallel direction, the conservation of momentum results in a decrease in the electron’s parallel momentum, while the conservation of angular momentum leads to an increase in the electron’s rotational energy, corresponding to the Anomalous Doppler Effect. It is important to note that there is no response during Cherenkov resonance for the electromagnetic wave, as the Cherenkov effect is primarily associated with electrostatic waves.

**IV. Critical Trapping Threshold of Anomalous Doppler Effect**

The Anomalous Doppler Effect functions as an effective damping force that impedes the electron acceleration process. By increasing the intensity of the electromagnetic wave, it is theoretically possible to balance the electrostatic field force, preventing further electron acceleration by the electrostatic field. The existence of this equilibrium will be demonstrated by varying the electromagnetic wave field intensity.

An electromagnetic wave with only a left-hand polarized circular component is considered, characterized by a wave vector and frequency , where , and k is aligned parallel to the static magnetic field. The electrostatic field and static magnetic field are set to and , respectively. As shown in Fig. 7, increasing the energy of the electromagnetic wave results in the parallel velocity becoming trapped in the resonant condition, ceasing to increase, while the perpendicular velocity continues to rise once the ratio exceeds a specific threshold. The electron’s orbit and momentum of the trapped electron are illustrated in Fig. 8.

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Figure 7. Time trace of velocity under different ratios of . (a) Vertical velocity, (b) Parallel velocity, (c) Zoom in parallel velocity.

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Figure 8. The electron’s orbit and momentum with trapped parallel energy

The threshold field can be determined by adjusting the electromagnetic wave intensity using a dichotomy control method, based on the final parallel velocity over a sufficiently long time. The critical ratios as functions of the dimensionless parameter are shown in Fig. 9, with the angles between and the wave vector k set at . The results indicate that a larger angle between k and leads to a lower critical ratio for trapping parallel energy. This reduction can be attributed to the increased electric field component along the parallel direction as the angle increases, enhancing the retarding force along that direction. Additionally, when the magnetic field strength BB and wave frequency are held constant, a decrease in the wave vector k results in a higher critical field strength. This is because a lower k corresponds to a higher resonant velocity, leading to an increase in the power imparted by the static electric field, . Consequently, a stronger electromagnetic wave intensity is required to achieve timely conversion of parallel energy.

With only a weak left-hand circularly polarized wave, it is possible to halt the increase in the electron's parallel momentum and transfer energy from the electrostatic field to rotational energy via the Anomalous Doppler Effect. For instance, in tokamaks, where the toroidal electric field is approximately , the threshold electric field for a left-hand circularly polarized wave to trap parallel energy is around in the plasma. This corresponds to an energy flux of approximately for , and with 0.02 T background magnetic field.

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Figure 9. The critical ratio of with the nomalized parameter . The , ,. The refractive index range is set from 4 to 50.

To validate the Anomalous Doppler Effect in high magnetic fields and assess its angle dependence, we consider a uniform magnetic field of and an electrostatic field , representative of typical tokamak startup conditions [21]. For a plane left-hand circularly polarized wave with parameters , , and , the wave’s energy flux is approximately . Utilizing the computational parallelism of a supercomputer, we simulate the interaction of 500 electrons with the wave at 500 distinct incident angles θ ranging from 0 to 90 degrees, thereby elucidating the angle dependence of the Anomalous Doppler Effect effect.

The simulation employs a time step of to ensure result convergence. The outcomes are depicted in Fig. 10, where the dashed yellow line indicates the earliest resonant time satisfying the condition . The results reveal that as θ\theta increases, the onset of resonance is delayed, and the velocity required for resonance becomes higher. Once resonance begins, the rotational velocity increases rapidly, while the parallel velocity becomes trapped in the resonant region, halting further growth.

A rainbow diagram of a rainbow

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Figure 10. Time evolution of (left) and (right) by electromgnetic wave with different wave incident angle .

**V. Electromagnetic wave drives the Anomalous Doppler Effect in magnetized plasma**

Runaway electrons pose a severe risk to tokamak devices, as their high-energy impacts can cause substantial damage to plasma-facing components (PFCs) [35]. Leveraging the Anomalous Doppler Effect within tokamaks provides a viable strategy for mitigating runaway electron energy, contingent upon satisfying three essential conditions. By fulfilling these criteria, it is possible to establish controlled resonant interactions, offering an effective means of suppressing runaway electrons.

1. The presence of a left-hand polarized wave component aligned with the electron motion in the magnetic field direction.
2. The phase velocity of the electromagnetic wave must remain subluminal, i.e., below the speed of light in a vacuum.
3. The wave must carry sufficient energy to counterbalance the accelerating velocity of the electrons.

The Appleton-Hartree formula is used to describe wave propagation in cold magnetized plasmas and is valid when the phase velocity significantly exceeds the thermal velocity of electrons. Considering the case where , we can analyze wave behavior under these conditions. The dispersion relation for a cold plasma in a magnetized medium is expressed as Eq. 19.

(19)

The is the wave frequency and k represents the wavevector. Both and depend on the electron cyclotron frequency(), the plasma frequency () and the ion cyclotron frequency(). Here, is the angle between and the static magnetic field . The dispersion relationship can be illustrated in Fig. 6(a), where the blue color represents the wavevector , while the red color represents 90 degrees. The polarization component in direction of the wave can be expressed as[1].

(20)

(21)

The electric field of the wave is expressed as . For a wave propagating along the zz-axis, the left-hand polarized wave is represented as , while the right-hand polarized wave is given by .

The left-hand and right-hand polarized wave component is caluculated with:

,

When B > 0 and , indicating that the wave is predominantly right-hand polarized. Conversely, the wave is primarily left-hand polarized.

The black dashed line in Fig. 11a denotes the boundary between the two polarization types, corresponding to in the dispersion relation for cold magnetized plasma. As shown in Fig. 11b, the region between the dashed lines predominantly features left-hand polarized waves, whereas outside these boundaries, the wave is mainly right-hand polarized.

In Fig. 11c, the black line represents waves in a vacuum. Only waves below this line, where the phase velocity , can drive the Anomalous Doppler Effect (ADE). Finally, as depicted in Fig. 11d, Anomalous Doppler Effect supporting waves are observed in both low- and high-frequency regions.

* In the low-frequency region, whistler waves and magnetized electron plasma waves are characterized by right-hand polarization.
* In the high-frequency region, extraordinary waves exhibit left-hand polarization when the angle θ is in close proximity to 90 degrees, as shown in Fig. 11b.

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Figure 11. Dispersion relationship in cold magnetized plasma. (a) Complete dispersion relationship in cold magnetized plasma. (b) Region dominated by left-hand polarization. (c) Region where the phase velocity is smaller than the speed of light in vacuum. (d) High-frequency and low-frequency regions.

(21)

When the runaway electron’s momentum satisfies the resonant condition (Eq. 21), it can excite intrinsic waves in the plasma. For simplification, we assume the velocity is aligned with the static magnetic field. In this analysis, we consider n=1 for the Anomalous Doppler Effect, and n=0 for Landau resonance. The case of n = −1 is excluded because the Normal Doppler Effect is not significant in the high-frequency region, as it requires a right-hand polarized wave. By combining Eq. 19 and Eq. 21, we derive the relationship between the wave properties and the resonant momentum, which is depicted in Fig. 12.

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Figure 12. In the high-frequency region, the lower-left boundary represents the upper hybrid wave at various angles. In the low-frequency region, the lower-right boundary corresponds to the lower hybrid wave at different angles. Panels (a) and (b) depict the dimensionless momentum for Landau resonance in the high-frequency region. Panels (c) and (d) illustrate the dimensionless momentum for the Anomalous Doppler Effect (ADE) in both high- and low-frequency regions. The unit for dimensionless momentum is expressed as .

In the high-frequency region for ADE, resonance curves with dimensionless resonance momentum greater than unity converge to the bottom-right region shown in Fig. 12c, which corresponds to the Extraordinary wave with a frequency range of . This explains the excitation of Extraordinary waves near these frequencies during runaway electron scattering in magnetized devices [28, 12]. Additionally, the dimensionless Landau resonant momentum in most of the low-frequency region is greater than 1, as shown in Fig. 12a, suggesting less wave attenuation by background thermal electrons, which facilitates wave formation in the high-frequency region. In the low-frequency region, when the energy of high-energy runaway electrons exceeds 10 MeV (with reduced momentum p > 20), the resonance curves of electromagnetic waves excited by the ADE effect typically pass through the top-left region depicted in Fig. 12d. This region is closely associated with the whistler wave zone, where whistler waves propagate parallel to the magnetic field. Thus, in Tokamak experiments, the observation of whistler waves is typically linked to the detection of high-energy electrons with energies exceeding 10 MeV [31]. In the low-frequency region, the dimensionless Landau resonance momentum is less than unity, as shown in Fig. 6b, indicating a higher degree of wave attenuation by background thermal electrons compared to the high-frequency region, making wave formation in this region more challenging.

Based on the above discussion, electromagnetic waves in the high-frequency region are more prone to exciting Anomalous Doppler Resonance due to polarization and damping effects, while waves in the low-frequency region are better suited for heating background electrons through Landau resonance, such as in lower-hybrid wave heating. Experiments have shown that runaway electrons can stimulate extraordinary waves with frequencies in the range to through ADE, thereby transferring their parallel energy to rotational energy [12, 29]. This suggests the potential for utilizing the reverse process — injecting extraordinary waves to suppress runaway electrons.

**VI. Launching Extraordinary Waves in Tokamaks for Runaway Electron Suppression**

The characteristic frequency with extraordinary mode in a tokamak plasma is shown in Fig. 13. Extraordinary waves slow down near the upper-hybrid frequency layer and reflect at the right-hand cut-off frequency layer. Therefore, it is necessary to inject the extraordinary wave from the high-field side of the tokamak, as it will be reflected at the right cut-off frequency when injected from the lower-field side. Different frequencies correspond to different upper-hybrid frequency layer positions, allowing the frequency of the extraordinary wave to be adjusted to align with regions where runaway events are more likely to occur, such as the core of the tokamak. Since the power requirement for trapping runaway electrons is on the order of watts, it is worth noting that precise frequency adjustments based on real-time plasma density diagnosis are possible to achieve, allowing for alignment with the tokamak core.

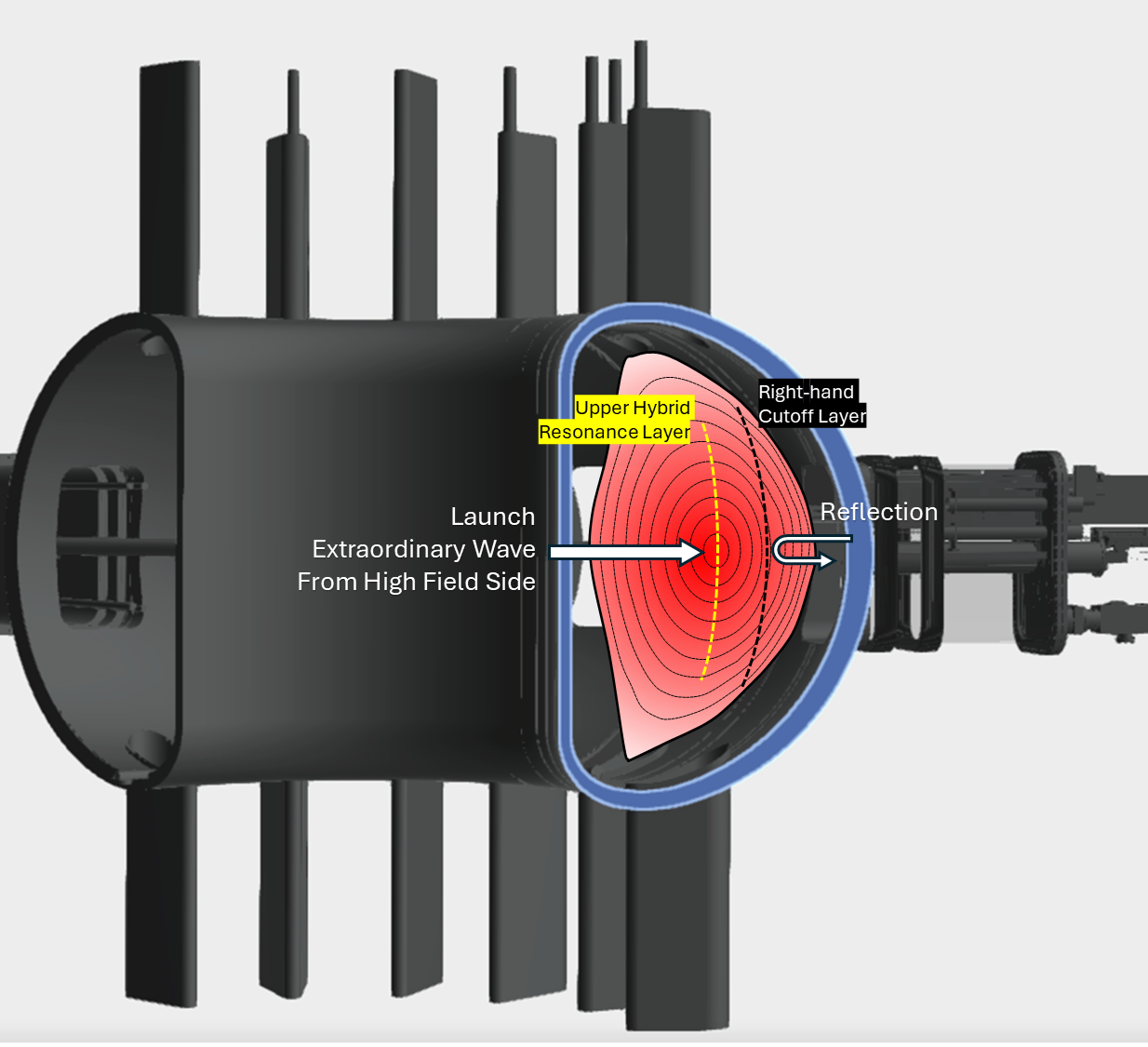


Figure 13. Characteristic frequencies in the tokamak. The Right-hand Cutoff layer is on the out side of the Upper Hybrid Resonance Layer .

**VII. Summary**

The Anomalous Doppler Effect has been identified as an effective mechanism for suppressing runaway electron energy by scattering the electron's parallel energy into perpendicular energy. Through studying the interaction between electrons and electromagnetic waves, this approach offers an innovative solution for runaway electron suppression in tokamaks, requiring a left-hand circularly polarized beam with an intensity of 9 W/m² based on the EAST startup scenario. In practical tokamak applications, the extraordinary wave predominantly contains left-hand circular polarization components and can be launched from the high-field side of the tokamak. Resonance at the upper hybrid layer drives the Anomalous Doppler Effect within the plasma core, effectively limiting the increase in runaway electron toroidal momentum. Numerical simulations demonstrate that when the electric field exceeds the critical threshold, the electromagnetic wave captures the parallel momentum of the electrons, continuously transferring energy from the parallel electrostatic field to rotational energy and resonant waves.

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